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# **Coupling between Power Systems and Pipelines during Geomagnetic Disturbances**

**L. Trichtchenko and D.H. Boteler**

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## **SUMMARY**

Design and operation of pipeline cathodic protection system for preventing pipeline steel corrosion as well as safety of crew during pipeline work need to take into account many factors. One of them is electromagnetic induction into a pipeline from the nearby power transmission lines.

The aim of the project is to examine the effects of electromagnetic induction from transmission line currents which includes not only 60 Hz currents, but also higher harmonics. These harmonics are produced by different technological and natural causes including geomagnetic storms.

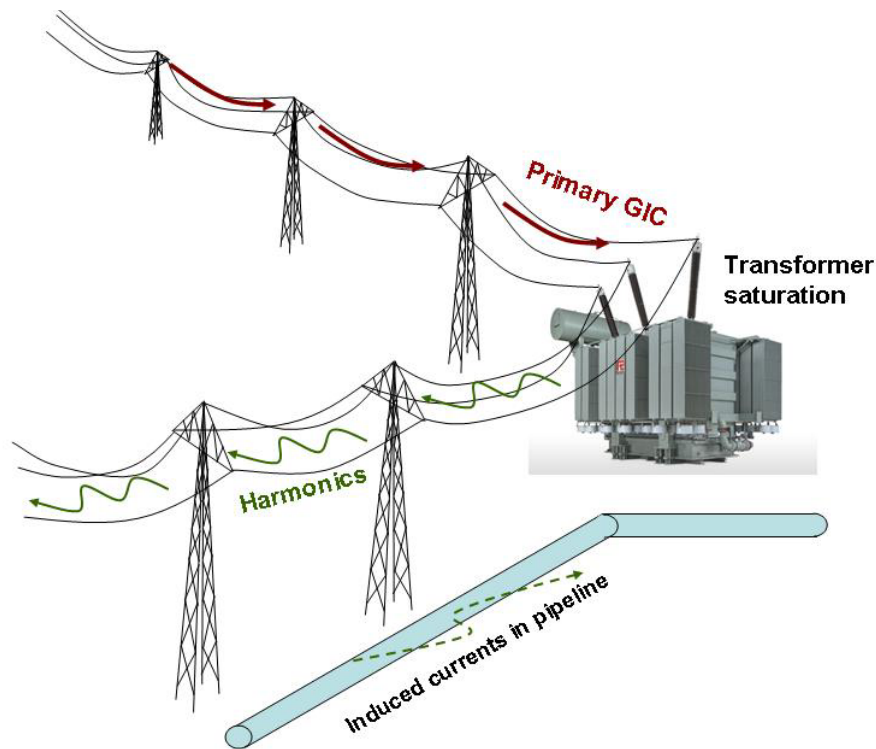
The results of the modelling show that fundamental and third harmonic produce equal amplitude of pipe-to-soil potential and current in the nearby pipeline. The simplified approach used in this “pilot” study clearly demonstrates the need for more substantial work with realistic pipeline and power lines network modelling as well as field measurements for model validation.

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## 1. Introduction

Disturbances of the geomagnetic field produced by space weather events cause variable geoelectric field at Earth's surface which drive electric currents ("geomagnetically induced" currents, GIC) in long conductors, such as power systems and pipelines. These "primary" GIC produce a variety of effects on both ground systems which were described in many papers, such as Molinksi (2002), Gummow (2002) and Trichtchenko and Boteler (2003). One of the consequences of the increased GIC in power system is the generation of harmonic currents due to saturation of transformers. These high frequency harmonic currents ( $N \cdot 60$  Hz, where  $N$  is an integer number) are capable of inducing "secondary" currents in adjacent pipelines (Figure 1). This report is concerned with the described aspect of "secondary" geomagnetic effects on pipeline operations, which has previously been overlooked.

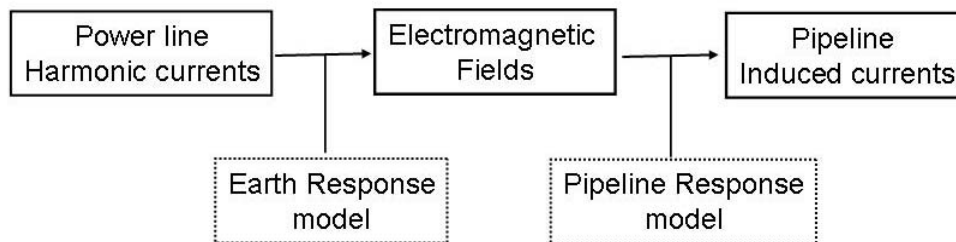


**Figure 1.** GIC in power systems cause transformer saturation, producing harmonic currents in power lines and induced currents in pipelines

The described process starts with GIC flowing through power transformers where the quasi-dc GIC creates an extra magnetic field in the transformer core. The resulting non-linear response of the transformer creates a distorted AC waveform that has a high content of harmonics.

Where pipelines share the same right-of-way with a power line, there is considerable electromagnetic coupling between the power line and the pipeline. There has been significant work on this issue (Taflove and Dabkowski, 1979; Dabkowski, 1999; Gummow et al, 1999), which has been almost exclusively concentrated on 60 Hz induction in pipelines. However, harmonics have also been observed in pipelines (Boteler et al. 2010; Boteler et al, 2013) and raised concerns about their contribution to the safe pipeline operation.

The results presented in this report are concerned with the impact of power system harmonics on pipe-to-soil potentials (PSP) and currents in pipelines sharing the same right-of-way with a power line. The report describes physical processes following the logic presented in Fig. 2.



**Figure 2.** Logic of the steps towards evaluation of the pipeline response to harmonic currents in power grids

Part 1 of the report describes the source and characteristics of the harmonic currents produced in power systems by GIC causing saturation of power transformers. The magnetic and electric fields produced by each harmonic at the location of the pipeline is calculated using the complex image method to account for the effects of induced currents in the Earth, as described in the Part 2.

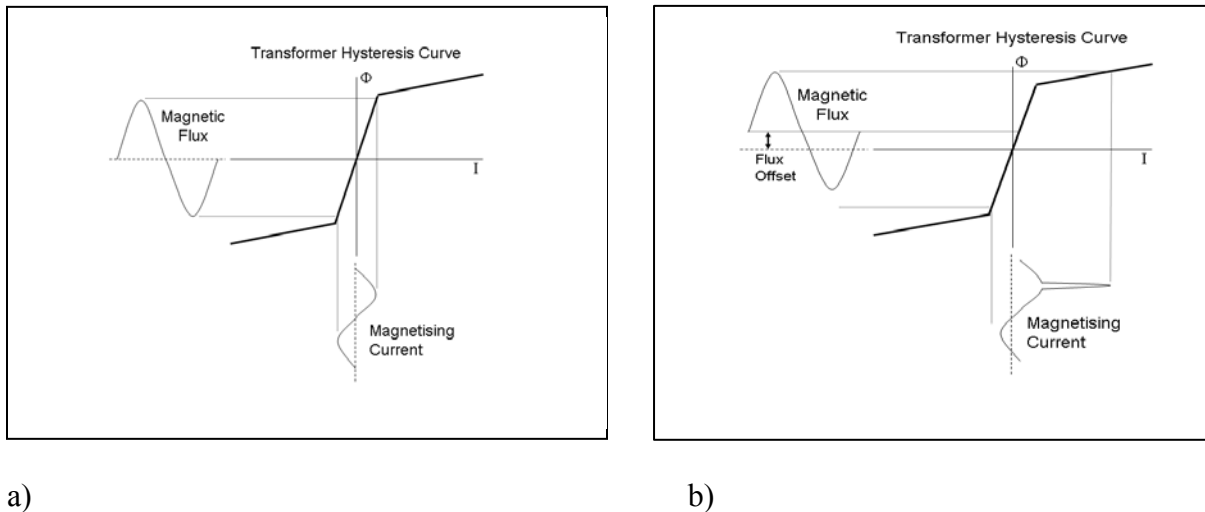
Part 3 shows how to set up the numerical model of the pipeline that includes frequency-dependent electromagnetic parameters of the pipeline circuit elements necessary for modelling the currents in the pipeline and the pipe-to-soil potential (PSP) produced by different harmonics.

Finally, in Part 4 all these components are combined to show the pipeline response resulting from the GIC-produced harmonics in the power system.

## 2. Harmonics in Power Systems

### 2.1. Cause of Harmonics

Geomagnetically induced currents (GIC) in power systems flow to ground through the windings of power transformers where these extra currents create partial saturation of the transformer core. This can be illustrated by considering the hysteresis curve of Figure 3.



**Figure 3.** Transformer hysteresis curves representing two cases:

- during normal operation;
- flux offset created by GIC shifts the operating regime into saturation

In normal operation (Fig. 3a), the sinusoidal magnetic flux is symmetrical about zero and is confined to the linear operating regime of the transformer. Accordingly, the flux produces a sinusoidal magnetising current which is symmetrical about zero.

When GIC flows through the neutral of transformer, the extra magnetic field produced shifts the operating point so that the peak of the AC cycle goes past the knee in the hysteresis curve producing a spike in the magnetising current which leads to the harmonics in this current. These AC harmonic currents (although much smaller than load current) flow along the transmission lines and create magnetic fields producing inductive coupling with nearby conducting lines (pipelines).

Fourier analysis of the distorted AC current waveform shows that, in addition to an increase in the fundamental (60Hz) component of the magnetising current, there are also components at the even and odd harmonics (120 Hz, 180 Hz etc.). Here the general formulas will be derived for the harmonic currents up to the third order.

## 2.2 Currents in the power line phase conductors

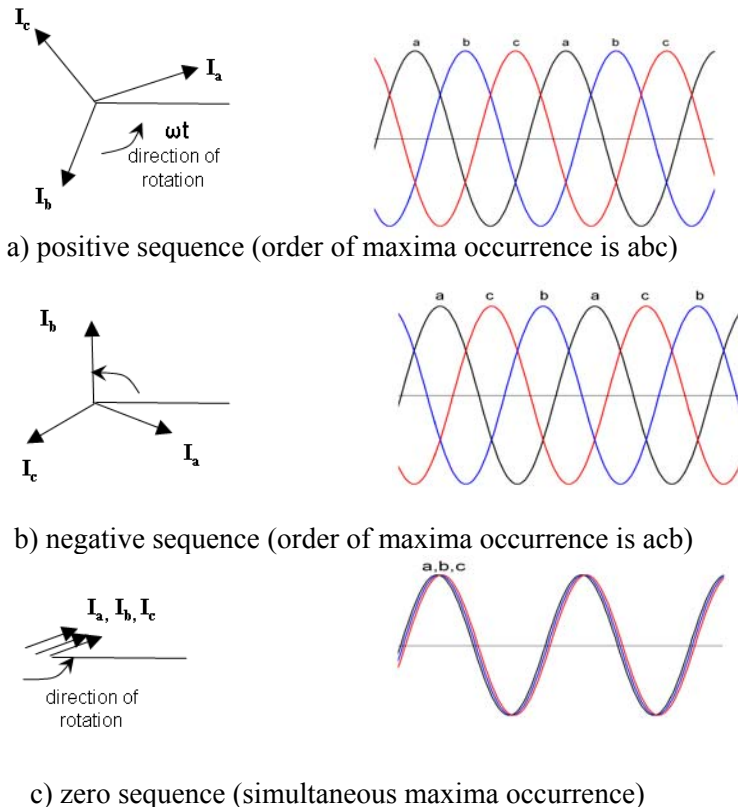
In three-phase operation, as used in most power systems, each of the three conductors in a transmission line is delivering the current from one phase of the three-phase generator. These three currents have the same amplitudes but the phases are 120 degrees apart (Equation (1)). Thus, in a well-balanced system the separate phase currents should add up to zero at the transformer neutral point.

The currents are expressed as follows:

$$\begin{aligned} i_a &= I_{1a} \cos(\omega t) \\ i_b &= I_{1b} \cos(\omega t + 240) \\ i_c &= I_{1c} \cos(\omega t + 120) \end{aligned} \quad (1)$$

where  $i_a$ ,  $i_b$ ,  $i_c$ , are the currents in each of three phase lines for the “fundamental” (i.e. 60Hz) frequency at each moment  $t$ ;  $I_{1a}$ ,  $I_{1b}$  and  $I_{1c}$  are their amplitudes, and  $\omega$  is angular frequency. The  $a$ ,  $b$ ,  $c$  subscripts are related to the sequence of the currents from the rotating machine as illustrated in Fig. 4a (so-called “positive sequence”).

When there is an imbalance in the system, these three phase currents do not add up to zero. In power system engineering this imbalanced current is usually represented as a combination of three specifically defined configurations of balanced “sequence components”. These are the positive sequence, as for the normal operations, negative and zero sequences, as shown in Figure 4.



**Figure 4.** Vector (left) and time-variation (right) schemes of the phase currents for three sequences used to present any unbalanced case: a) positive, b) negative, and c) zero



The concept of “sequence components” was first proposed in 1895 and later the mathematical analysis was done in 1912 and 1915 and since then it has been widely used in power engineering (e.g. Wagner and Evans, 1933; Blackburn, 1993 and references therein).

When higher harmonics are presented, currents in each phase (a, b and c) can be represented as Fourier series, i.e.

$$\begin{aligned} i_a &= I_{1a} \cos(\omega t) + I_{2a} \cos(2\omega t) + I_{3a} \cos(3\omega t) + \dots \\ i_b &= I_{1b} \cos(\omega t + 240) + I_{2b} \cos(2(\omega t + 240)) + I_{3b} \cos(3(\omega t + 240)) + \dots \\ i_c &= I_{1c} \cos(\omega t + 120) + I_{2c} \cos(2(\omega t + 120)) + I_{3c} \cos(3(\omega t + 120)) + \dots \end{aligned}$$

where index 1 is for fundamental (60Hz), 2 for the second harmonic (120 Hz) and 3 for the third harmonic (180 Hz).

After re-arrangements and subtraction of any multiple of 360 in *b* and *c* phases, we have

$$\begin{aligned} i_a &= I_{1a} \cos(\omega t) + I_{2a} \cos(2\omega t) + I_{3a} \cos(3\omega t) + \dots \\ i_b &= I_{1b} \cos(\omega t + 240) + I_{2b} \cos(2\omega t + 120) + I_{3b} \cos(3\omega t) + \dots \\ i_c &= I_{1c} \cos(\omega t + 120) + I_{2c} \cos(2\omega t + 240) + I_{3c} \cos(3\omega t) + \dots \end{aligned} \quad (2)$$

Thus, the harmonic components in each phase conductor *a*, *b* and *c* have phase relations as shown in Table 1.

**Table 1.** Phase relations of harmonics on a 3-phase system

Harmonic	Frequency	Phase of $i_a$	Phase of $i_b$	Phase of $i_c$	Phase Sequence
1(fundamental)	60Hz	0	240	120	Positive
2	120Hz	0	120	240	Negative
3	180Hz	0	0	0	Zero
4	240Hz	0	240	120	Positive
5	300Hz	0	120	240	Negative
6	360Hz	0	0	0	Zero

From these equations it follows that summations of the harmonic currents from all three phases reduces any frequency except multiple of 3, because the 3-rd harmonic has zero phase sequence. Therefore, currents induced in the pipeline might have significant response to the third harmonic.

### 3. Electric Fields at the Earth surface produced by the currents in the power lines

The harmonically-varying currents in the power lines create a changing magnetic field and electric fields in the surrounding media and at the pipeline location. Assuming the amplitudes of fundamental and harmonic currents are monitored in the power grid, our next task (following the logic presented in Fig. 2) is to find the electric field in the earth produced by these currents.

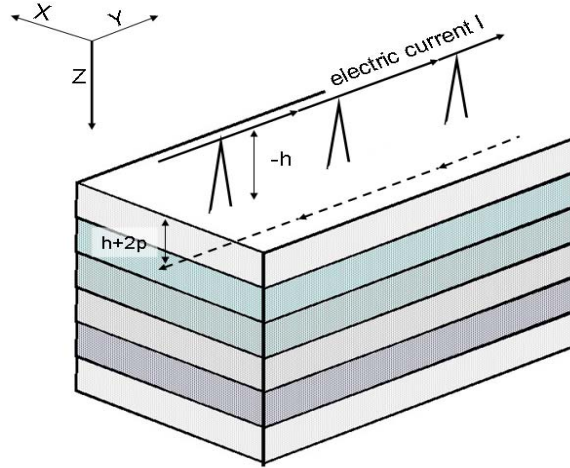
The electric field  $E$  produced at the Earth's surface by an external source satisfies the diffusion equation in the form:

$$\nabla^2 E = i\omega\mu\sigma E \quad (3)$$

where  $i = \sqrt{-1}$ ,  $\omega$  is angular frequency of the source,  $\mu$  is magnetic permeability and  $\sigma$  is conductivity of the media.

#### 3.1. Electric field calculations for a line current source

The geomagnetic coordinate system, i.e. X-axis is directed northward, Y-axis is directed eastward and Z-axis downward (Fig. 5), with line current  $I$  located at the height of  $-h$  (i.e. current in the overhead power line) and is directed along Y-coordinate has been used.



**Figure 5.** Schematic representation of a current in the overhead wire (at a height  $-h$ ) and image current in the Earth (at a depth  $h+2p$ )

The solution of the equation (3) at the surface of the Earth can be written as (Hermance and Peltier, 1970; Wait, 1992; Boteler and Pirjola, 1998)

$$E_y = -I \frac{i\omega\mu_0}{2\pi} \int_0^\infty (1 - R(\nu)) e^{-h\nu} \frac{\cos \nu x}{\nu} d\nu \quad (4)$$

where  $R$  is the reflection coefficient representing properties of the conducting earth,  $\nu$  is the integration variable called “wavenumber” representing the geometrical properties of the

cylindrical electromagnetic wave created by a line current. The result of integration will be independent of this variable.

The reflection coefficient can be expressed as

$$R = \frac{i\omega\mu_0/Z - \nu}{i\omega\mu_0/Z + \nu}, \quad (5)$$

where  $Z$  is the surface impedance (see the paragraph 3.2 below).

To solve (4) analytically, the reflection coefficient can be approximated as (Boteler and Pirjola, 1998)

$$R = \frac{1/p - \nu}{1/p + \nu} \approx \exp(-2p\nu) \quad (6)$$

where

$$p = \frac{Z}{i\omega\mu_0} \quad (7)$$

Substituting for  $R$  in equation (4) gives the expression for the electric fields:

$$E_y = -I \frac{i\omega\mu_0}{2\pi} \ln \frac{\sqrt{(h+2p)^2 + x^2}}{\sqrt{h^2 + x^2}} \quad (8)$$

This shows that the electric field at the surface of the Earth can be thought of as the sum of an external part (produced by a line source) and an internal part (produced by an image current at a “complex” depth  $h+2p$ ).

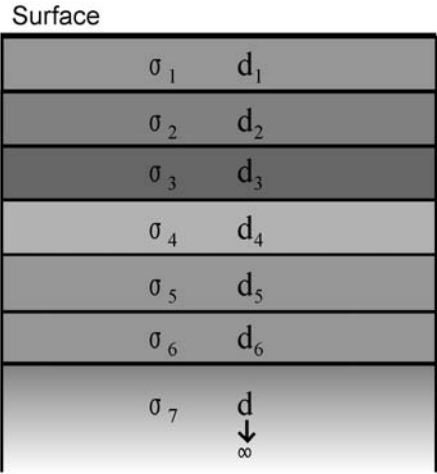
Equation (4) or in more simple form Eq. (8) together with either (5) or (6) and (7) can be used to calculate the approximate values of electric fields at the Earth's surface at location  $x$  from the vertical projection of phase current  $I$  to the earth surface.

### 3.2. Surface Impedance for a Layered Earth model

For the “layered” earth model (Figure 6), the unknown impedance  $Z$  can be found from general formulas of impedance at the surface of any layer  $n$  (where  $n=1$  is the surface of the Earth and  $n=N$  is the last layer before the uniform half-space) by applying the recursion relation (Weaver, 1994).

$$Z_n = i\omega\mu \left( \frac{1 - r_n e^{-2k_n d_n}}{k_n (1 + r_n e^{-2k_n d_n})} \right) \quad (9)$$

where  $d_n$  and  $k_n$  are thickness and propagation constant of layer  $n$ , with propagation constant depending on the layer's conductivity as:  $k_n = \sqrt{i\omega\mu_0\sigma_n}$  and  $r_n$  are the reflection coefficients at each layer boundary (except the last layer  $N$ )



**Figure 6.** Schematic representation of the layered earth model with conductivities  $\sigma_n$  and thicknesses  $d_n$  of each layer.

The reflection coefficients are defined as:

$$r_n = \frac{1 - k_n \frac{Z_{n+1}}{i\omega\mu}}{1 + k_n \frac{Z_{n+1}}{i\omega\mu}} \tag{10}$$

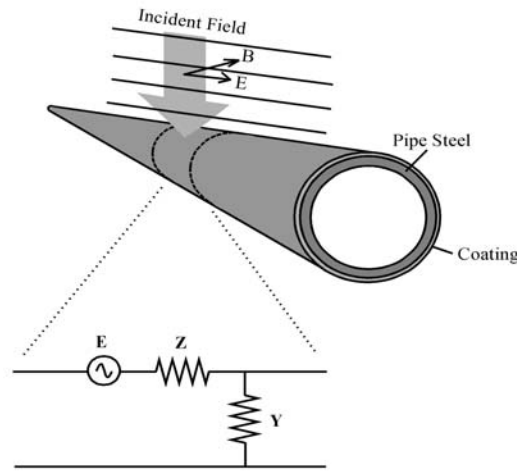
and for the last layer

$$Z_N = \frac{i\omega\mu}{k_N} \tag{11}$$

## 4. Pipeline Response to Harmonics

### 4.1. Network model of a pipeline

The electric field induced in the pipeline produces variations in the pipe-to-soil potential (PSP). The exact PSP produced depends on the electrical characteristics of the pipeline. This can be modelled using transmission line theory (Taflove and Dabkowski, 1979). The starting point is the approximation of pipeline network by the series of uniform sections with the geoelectric field  $E$  in each section (Figure 7).



**Figure 7.** Schematic representation of the uniform element of a pipeline network.

The voltage and current along the line are given by

$$\frac{dV}{dx} + ZI = E \quad (11)$$

and

$$\frac{dI}{dx} + YV = 0 \quad (12)$$

from which differentiation and substitution lead to the equations

$$\frac{d^2V}{dx^2} - \gamma^2 V = \frac{dE}{dx} \quad (13)$$

and

$$\frac{d^2I}{dx^2} - \gamma^2 I = -YE \quad (14)$$

where  $Z$  and  $Y$  are the series impedance along the pipeline and the parallel admittance of the pipeline coating. These are used to determine the characteristic impedance,  $Z_0$ , and the propagation constant,  $\gamma$ , given by

$$Z_0 = \sqrt{\frac{Z}{Y}} \quad \gamma = \sqrt{ZY}$$

If the electric field is assumed to be uniform within a section of pipeline,  $dE/dx = 0$  and equation (13) becomes

$$\frac{d^2V}{dx^2} - \gamma^2V = 0 \quad (15)$$

Equations (14) and (15) have solutions of the form

$$V = Ae^{\gamma x} + Be^{-\gamma x} \quad (16)$$

and

$$I = \frac{E}{\gamma Z_0} - \frac{A}{Z_0} e^{\gamma x} + \frac{B}{Z_0} e^{-\gamma x} \quad (17)$$

where  $A$  and  $B$  can be found from conditions at the ends of the pipeline.

At  $x=0$  equation (16) gives

$$V_i = A + B \quad (18)$$

While at the end of the section,  $x = L$ , equation (16) gives

$$V_k = Ae^{\gamma L} + Be^{-\gamma L} \quad (19)$$

Combining these two equations gives

$$A = \frac{V_k - V_i e^{-\gamma L}}{e^{\gamma L} - e^{-\gamma L}} \quad (20)$$

and

$$B = \frac{V_i e^{\gamma L} - V_k}{e^{\gamma L} - e^{-\gamma L}} \quad (21)$$

Substituting for  $A$  and  $B$  in equation (16) for voltage gives

$$V = \left( \frac{V_k e^{\gamma L} - V_i}{e^{\gamma L} - e^{-\gamma L}} \right) e^{-\gamma(L-x)} + \left( \frac{V_i e^{\gamma L} - V_k}{e^{\gamma L} - e^{-\gamma L}} \right) e^{-\gamma x} \quad (22)$$

Similarly substituting for  $A$  and  $B$  in equation (17) for current gives

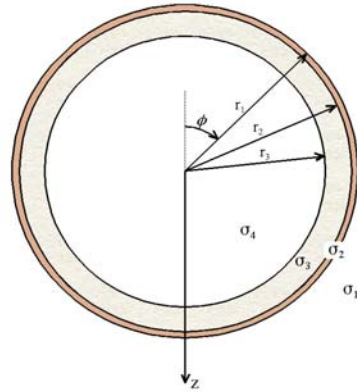
$$I = -\frac{1}{Z_0} \left( \frac{V_k e^{\gamma L} - V_i}{e^{\gamma L} - e^{-\gamma L}} \right) e^{-\gamma(L-x)} + \frac{1}{Z_0} \left( \frac{V_i e^{\gamma L} - V_k}{e^{\gamma L} - e^{-\gamma L}} \right) e^{-\gamma x} + \frac{E}{\gamma Z_0} \quad (23)$$

This represents the variation of voltage and current within a uniform pipeline section. Combining these sections as Thevenin equivalent circuits (Taflove and Dabkovski, 1979) or pi-circuits (Boteler, 2013), induced currents and voltages can be calculated.

## 4.2. Frequency dependent pipeline parameters

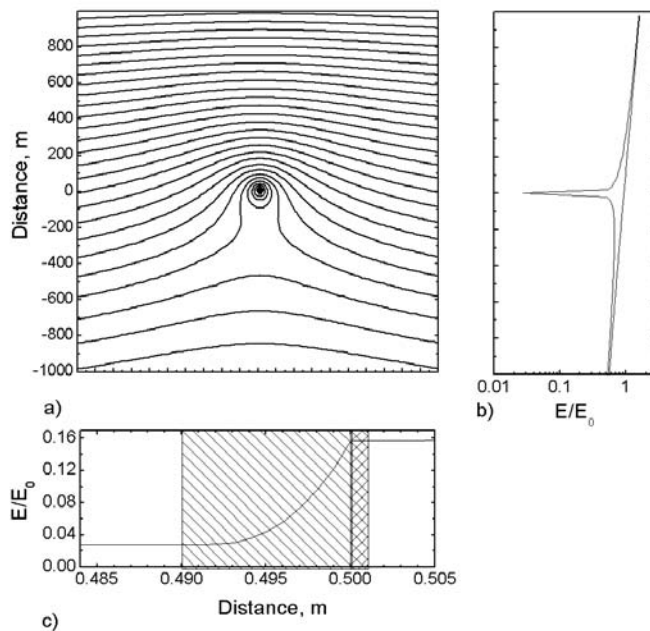
The key characteristics of the pipeline are the series impedance,  $Z$ , along the pipeline and the parallel admittance,  $Y$ , of the pipeline coating. Assuming the pipeline has a simple multilayered structure as shown in Figure 8 with layers 1,2,3,4, are soil, coating, steel and

gas (oil) respectively with their electromagnetic parameters,  $\sigma$ , conductivity,  $\epsilon$ , dielectric permittivity and  $\mu$ , relative magnetic permeability.



**Figure 8.** Schematic representation of the pipeline as a layered cylinder

The response of the pipeline to the external electromagnetic wave (with electric field parallel to the pipeline axis) can be modelled as diffusion equation in cylindrical coordinate system. The exact solution of this equation with appropriate boundary conditions is presented by Trichtchenko and Boteler (2001). Example of the variations of electric field outside and inside the pipeline layers (normalised to the electric field in the absence of the pipeline) for 60 Hz frequency is shown in Figure 9.



**Figure 9.** Example of the modelling of the layered cylinder. The normalized absolute value of the electric field,  $f=60$  Hz,

- a) contour lines; b) values along cross section at  $X=0$  and  $X=1000$ m
- c) magnified view across the pipeline wall

Thus, for frequencies 60 Hz and higher the electric field does not uniformly penetrate the steel and the depth of penetration decreases with increasing frequency. Therefore the impedance along the pipeline is frequency dependent.

Total impedance is the sum of the external and internal impedances.

$$Z_{total} = Z_{ext} + Z_{int} \quad (24)$$

In this formula the total impedance represents the ratio of the current in the pipeline steel to the electric field in the absence of the pipeline (i.e. induced in the ground).

For the external impedance, the simple representation, so called Dubanton equation (Gary, 1976; Deri et al, 1981) exists which is a good approximation of the Carson (1926) formulas. This simple representation is used also by Dabkowski (1999) for calculations of the longitudinal electric field in pipelines produced by the power line interference.

$$Z_{ext} = \frac{i\omega\mu_1}{2\pi} \ln \frac{2}{k_1 R_2} \quad (25)$$

where index 1 is referred to the surrounding medium (soil), and  $R_2$  is the outer radius of the pipeline steel .

For the internal impedance (with external return), which represents the AC-resistance of the pipeline steel, we use approximation given by Sunde (1968), his equation (1.111).

$$Z_{int} = \frac{2\gamma\delta\rho}{[\exp(\gamma\delta) - \exp(-\gamma\delta)]} = \frac{\gamma\delta\rho}{sh(\gamma\delta)} \quad (26)$$

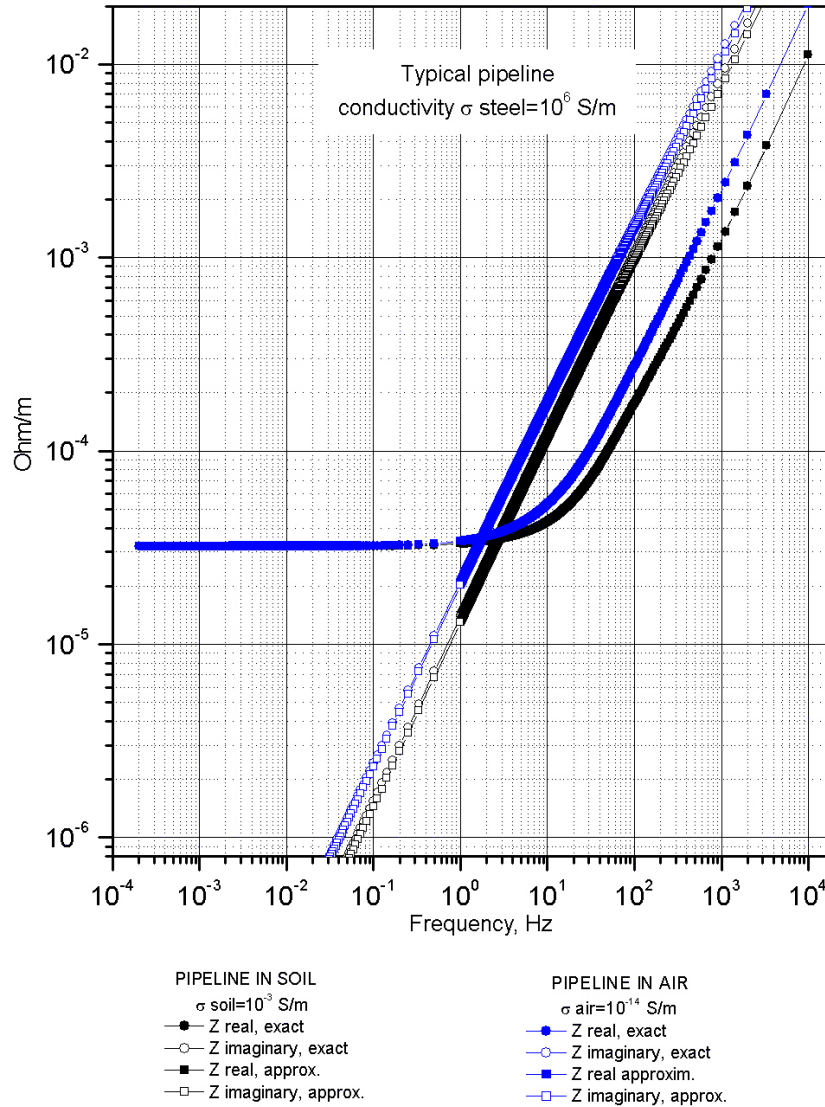
where

$$\gamma = \sqrt{\frac{i\omega\mu_3\rho}{2\pi\delta R_{av}}} \quad (27)$$

$R_{av}=(R_2+R_3)/2$  is mean radius of the pipeline steel layer,  $\delta$  is thickness of the steel and  $\rho$  is DC-resistance per unit length,  $\mu_3$  is steel permeability, and  $\rho = 1/(\pi\sigma_3(R_2^2 - R_3^2))$  where index 3 refers to the internal surface of the pipeline steel layer, index 2 to the external surface of the steel layer and index 1 to the external radius of the pipeline together with coating.

Verification of the above expressions for the frequency dependent impedances has been carried out by Trichtchenko et al, (2004) and is presented in Figure 10.





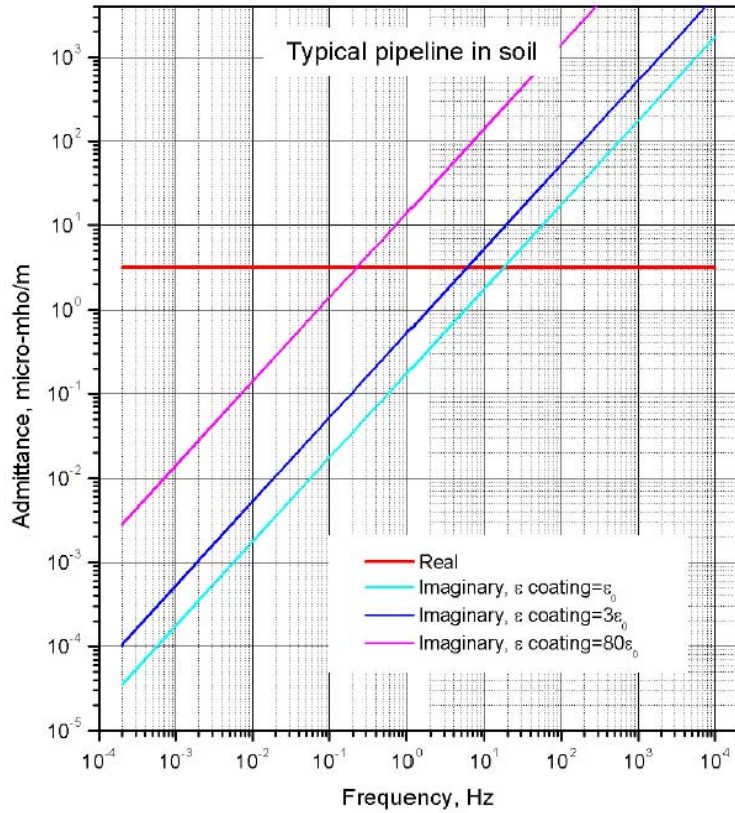
**Figure 10.** Comparisons of the impedances calculated by exact method (Trichtchenko et al, 2004) and analytical approximations (equations 24-26) for pipeline in soil and in air.

To complete the derivation of the frequency-dependent parameters of the pipeline, we also need to know the admittance  $Y$  through the pipeline coating to ground. This is a complex value and includes both the capacitance  $C$  between pipeline steel as one plate and surface of the earth as another as well as conductance through the coating. Therefore, the admittance is

$$Y = Q + i\omega C \quad (28)$$

where  $Q=2\pi\xi_2R_l$  with  $R_l$  is external radius of pipeline (with coating),  $\xi_2$  is coating conductance ( $S/m^2$ ), and  $C=2\pi\varepsilon_2/l\ln(R_1/R_2)$ ,  $\varepsilon_2$  is coating permittivity.

The variation of the admittance with frequency is shown below in Figure 11.



**Figure 11.** Real and imaginary parts of a complex admittance of pipeline for 3 cases of different coating permittivities and external radius 0.457 m, internal radius is 0.448 m, coating conductance is  $2 \cdot 10^{-5} S/m^2$ , thickness 0.00025 m

## 5. Effects of harmonics on pipeline cathodic protection

The previous sections have shown how the frequency and phase relations of harmonic currents affect the electric field induced in the pipeline; how the complex image method can be used to account for the effects of induced currents in the Earth; and how the pipeline response to a specified electric field can be modelled using DSTL theory and the NPM method. Now we combine all these parts to calculate the PSP produced in a pipeline by a realistic set of harmonic currents flowing in an adjacent power line.

The Steps to calculate the currents and voltages in the pipeline induced by the harmonics are:

- i) Specify a spectrum of harmonic currents
- ii) Calculate the electric fields produced by each harmonic
- iii) Calculate the PSP produced by each harmonic electric field

### 5.1. Spectrum of harmonic currents

In our example only fundamental, second and third harmonics will be discussed. The fundamental corresponds to a positive sequence, second to a negative sequence, third to a zero sequence, i.e. by this example all three sequence configurations are covered.

Fundamental current is 1000A, second and third are 1%, which is 10A each (power engineer, personal communication, 2013).

### 5.2. Ground electric field produced by each harmonic

To illustrate the calculations of electric field produced by each harmonic at the Earth surface, consider the three-phase transmission line configuration as shown in Figure 12.

Electric field produced by each harmonic current can be presented as:

Fundamental (positive sequence):

$$E_1 = I_1 \left[ (E_{1aL} + E_{1aR}) + (E_{1bL} + E_{1bR}) \cos(240^\circ) + (E_{1cL} + E_{1cR}) \cos(120^\circ) \right] \quad (29)$$

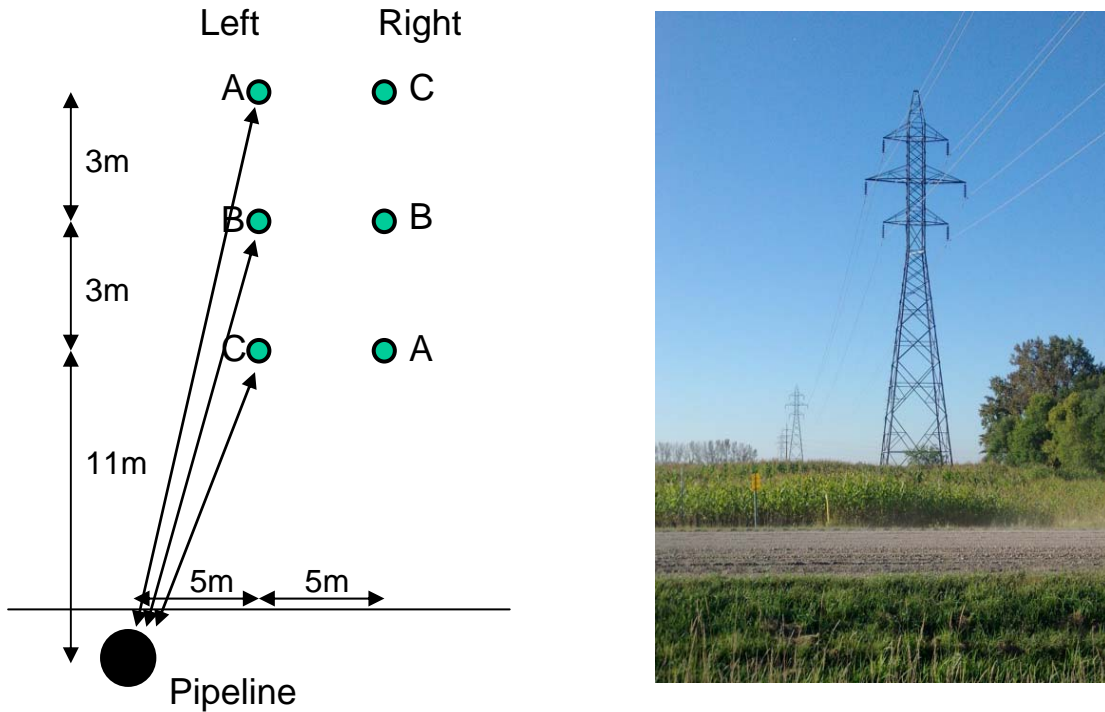
Second Harmonic (negative sequence):

$$E_2 = I_2 \left[ (E_{2aL} + E_{2aR}) + (E_{2bL} + E_{2bR}) \cos(120^\circ) + (E_{2cL} + E_{2cR}) \cos(240^\circ) \right] \quad (30)$$

Third Harmonic (zero sequence):

$$E_3 = I_3 \left[ (E_{3aL} + E_{3aR}) + (E_{3bL} + E_{3bR}) + (E_{3cL} + E_{3cR}) \right] \quad (31)$$

where  $E_{1aL}$ ,  $E_{1aR}$ ,  $E_{2aL}$ ,  $E_{2aR}$ ,  $E_{3aL}$ ,  $E_{3aR}$  are normalised electric fields (i.e. electric field per unit current) produced by the current in the phase  $a$  at the frequency of each harmonic (1,2 or 3) and by each conductor (left and right) respectively. Similar notations are used for normalised electric fields produced by phase  $b$  and phase  $c$ .



**Figure 12.** Arrangement of power line/pipeline used for calculations; a) Diagram showing distances from conductors to pipeline; b) Photo of the transmission line next to a pipeline (yellow marker)

The normalised electric fields at the earth surface produced by each separate line at each of three frequency can be found from (8), i.e.

$$E/I = -\frac{i\omega\mu_0}{2\pi} \ln \frac{\sqrt{(h+2p)^2 + x^2}}{\sqrt{h^2 + x^2}} \quad (32)$$

where  $h$  is the height of each conductor above the surface,  $x$  is the horizontal distance from the conductor to pipeline location (as per Figure 4a) and  $h+2p$  is the complex depth location of the image current with respect to the earth surface (Eq. (7)).

For our calculations the layered earth model has been chosen to have 11 layers and two different arrangements, such as

Case 1, “Resistive Earth” model with depths (in meters) as:

100, 1000, 10000, 14000, 75000, 100000, 160000, 110000, 150000, 230000,  $10^9$

and conductivities (in 1/Ohm-m) as:

0.01, 0.0077, 0.0002, 0.008, 0.0033, 0.0063, 0.0346, 0.1258, 0.4168, 1.122, 2.0892

Case 2, the hypothetical layered earth with top sea water layer with same depths of layers but conductivities (in 1/Ohm-m) as:

1.00, 0.0077, 0.0002, 0.008, 0.0033, 0.0063, 0.0346, 0.1258, 0.4168, 1.122, 2.0892

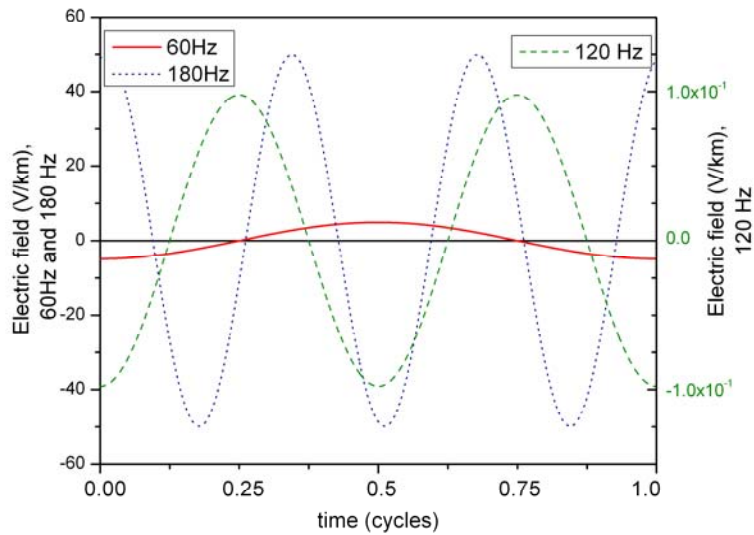
Results of these calculations are presented in Table 2 (case 1) and Table 3 (case 2) and Figures 13 and 14.

**Table 2.** results of electric field calculations for Case 1 (Resistive Earth) model

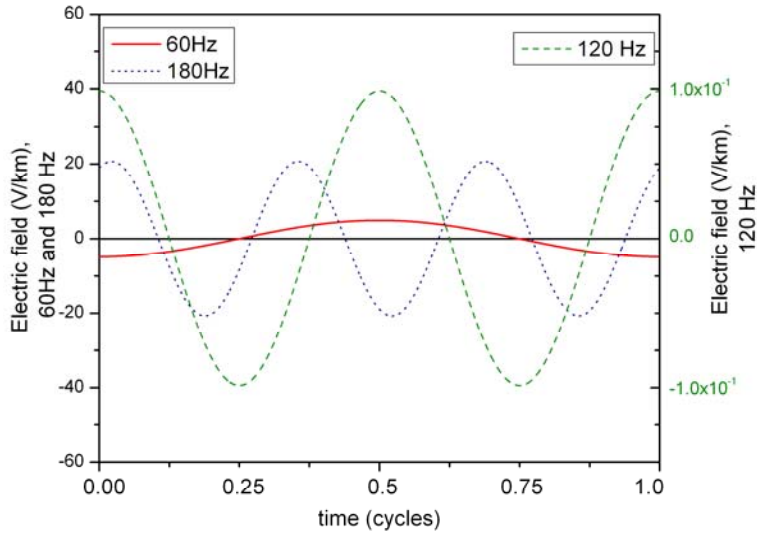
Freq	Surface Impedance (amplitude)	Surface Impedance (phase)	Electric Field (amplitude)	Electric Field (phase)
60 Hz	0.222 Ohm	42.811 degrees	$4.8910^{-3}$ V/m	-90 degrees
120 Hz	0.329 Ohm	44.008 degrees	$9.779 \cdot 10^{-5}$ V/m	-90 degrees
180 Hz	0.404 Ohm	43.163 degrees	$4.998 \cdot 10^{-2}$ V/m	77.46 degrees

**Table 3.** results of electric field calculations for Case 2 (Seawater top) model

Freq	Surface Impedance (amplitude)	Surface Impedance (phase)	Electric Field (amplitude)	Electric Field (phase)
60 Hz	$2.017 \cdot 10^{-2}$ Ohm	44.758 degrees	$4.911 \cdot 10^{-3}$ V/m	-90 degrees
120 Hz	$3.055 \cdot 10^{-2}$ Ohm	46.151 degrees	$9.855 \cdot 10^{-5}$ V/m	90 degrees
180 Hz	$3.788 \cdot 10^{-2}$ Ohm	45.379 degrees	$2.072 \cdot 10^{-2}$ V/m	66.122 degrees



**Figure 13.** Results of the electric field calculations for Case 1, Resistive Earth. The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz). Note the scales for 2<sup>nd</sup> harmonic–associated electric field (green lines) are different (shown in green).



**Figure 14.** Results of the electric field calculations for Case 2 (Seawater top) layered earth models. The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz). Note the scales for 2<sup>nd</sup> harmonic-associated electric field (green lines) are different.

### 5.3. Harmonic Pipe-to soil potential and currents in pipeline

To illustrate the calculations of PSP and currents produced by harmonics, an example pipeline with the following parameters was used.

Uniform pipeline section (no bends) with the length of 450 km is running along the power line with terminating impedances of 0.1 Ohms at both ends of the pipeline. The low values of terminating impedances are chosen in order to anticipate possible continuations of pipeline at both ends. The pipe-to soil potential and current are calculated at the distance of the 185 km from the beginning of the pipeline.

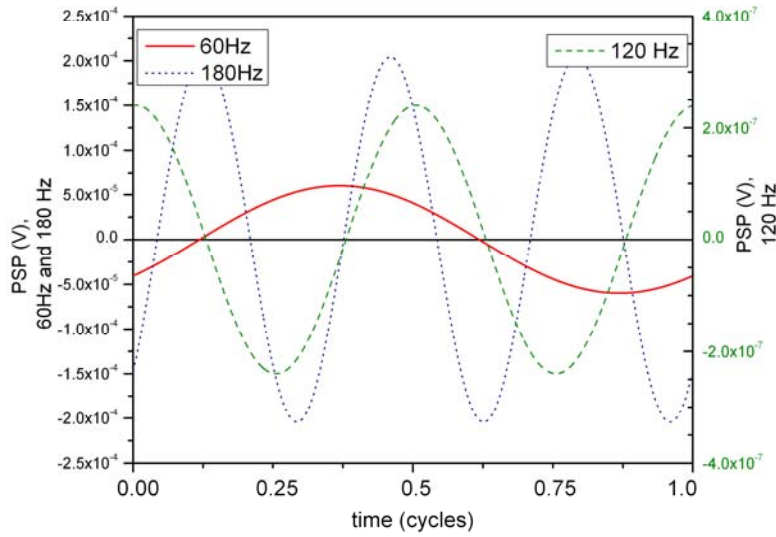
Table 4, Geometric and electromagnetic characteristics of pipeline

Dimensions of pipeline		Electromagnetic Properties	
External Radius	0.457 m	Steel Conductivity	$5 \cdot 10^6$ S/m
Internal Radius	0.448 m	Coating Conductance	$20 \cdot 10^{-6}$ S/m <sup>2</sup>
Coating Thickness	0.00025 m	Steel Permeability	150 (relative)
		Coating Permittivity	3 (relative)

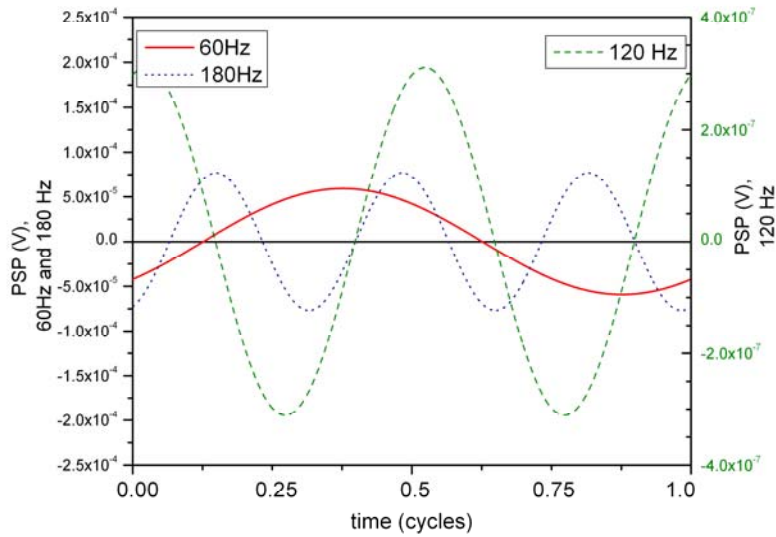
Table 5, Pipeline impedances and admittances at 3 different frequencies

Freq Hz	Pipeline Impedance (real) Ohm/km	Pipeline Impedance (imaginary) Ohm/km	Coating Admittance (real) S/km	Coating Admittance (imag) S/km
60	0.089	0.572	0.0574	0.115
120	0.160	1.074	0.0574	0.230
180	0.229	1.554	0.0574	0.344

Variations of the pipe-to-soil potential and current along the pipeline for two cases surface impedance (or the layered earth models) are presented in the following figures.

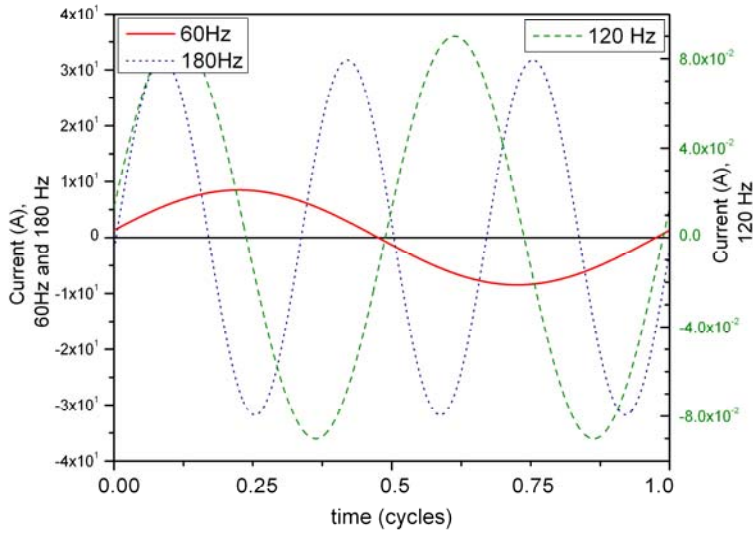


**Figure 15** Variations of PSP at Km 185 due to harmonics with a resistive earth model (case 1). The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz)

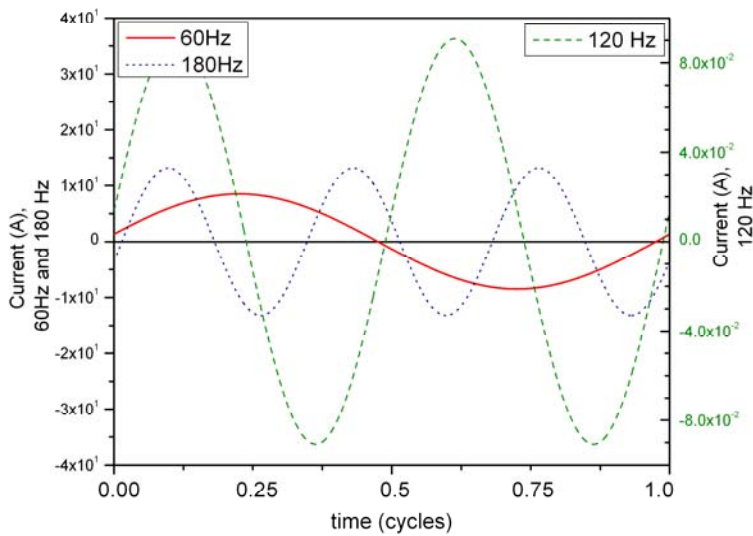


**Figure 16.** Variations of PSP at Km 185 due to harmonics with a top sea water model (case 2). The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz)





**Figure 17** Variations of pipeline current due to harmonics with a resistive earth model (case 1). The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz)



**Figure 18.** Variations of pipeline current due to harmonics with a top sea water model (case 2). The currents in power transmission lines are set to 1000A (60Hz), 10A (120Hz), 10A (180Hz)



## **6. Conclusions**

This report examines the effects on pipelines from harmonic currents produced in power lines during geomagnetic disturbances. On 3-phase lines the different harmonics have different phase relations and are classified as ‘positive sequence’, ‘negative sequence’, and ‘zero sequence’.

The transmission line currents can be used with an earth conductivity model to calculate the electric fields produced at the locations of adjacent pipelines. Pipeline network modelling can then be used to determine the pipe-to-soil potential (PSP) variations and pipeline currents that are produced.

All these steps can be combined to calculate the PSP and pipeline currents produced by harmonic currents in a nearby power line.

The results show that PSP and current in pipeline due to the third harmonic can be higher than those due to the unbalanced fundamental frequency in case of the resistive top layer of the layered earth model and comparable with those induced by fundamental frequency if the top layer is conductive (sea water).

Further work is needed to validate the calculations by more experimental measurements in the field. As well, application to more realistic arrangements of pipeline and transmission line networks configurations needs to be examined

## 7. Appendix

### 7.1 FORTRAN Code 1

Calculation of electric field in the ground induced by power line harmonic currents  
Earth conductivity model as per Ontario 2013 report,  
ON1-zone 1 of the Ontario report (Case 1 of this report). ON2-top layer of Zone 1 changed  
to sea water (1.0 Ohm-m), case 2 of this report  
Bottom layer has been changed to thickness of 1E9m  
Ea,b,c-electric field produced by each phase, left and right sides of tower, normalised by  
amplitude current,  
P is skin-depth (dependent on freq), h and x are locations with respect to pipeline

```
parameter (npoint=3) !number of harmonics (npoint=1 is 60 Hz)
REAL dOn1(11),sigmaOn1(11),dOn2(11),sigmaOn2(11), Cur(npoint)
REAL ZR(npoint),ZI(npoint),ZA(npoint), Zph(npoint), F(npoint)
REAL AMPLE(npoint),Phase(npoint)
COMPLEX cOn1(11),cOn2(11)
COMPLEX Z(npoint),Z1(npoint),ZZ(npoint),E(npoint)
Complex EaL(npoint), EbL(npoint), EcL(npoint), P(npoint)
Complex EaR(npoint), EbR(npoint), EcR(npoint)
```

C Power Tower Arrangements in meters from pipeline

```
XcL=5.
XaL=5.
XbL=5.
HcL=11.
HbL=14.
HaL=17.
```

```
XcR=10.
XbR=10.
XaR=10.
HcR=17.
HbR=14.
HaR=11.
```

C currents for each harmonic

```
Cur(1)=1000.
Cur(2)=10.
Cur(3)=10.
```

C Ontario model

```
data dOn1/100.,1000.,10E3,14E3,75E3,100E3,160E3,110E3
* ,150E3,230E3,1E9/
data sigmaOn1/0.01,0.0077,0.0002,0.008,0.0033,0.0063,0.0346,0.1258
* ,0.4168,1.122,2.0892/
```

```

C Ontario model with top water
  data dOn2/100.,1000.,10E3,14E3,75E3,100E3,160E3,110E3
  *,150E3,230E3,1E9/
  data sigmaOn2/1.,0.0077,0.0002,0.008,0.0033,0.0063,0.0346,0.1258
  *,0.4168,1.122,2.0892/

  NON=11
  PI=3.1415926
  TOPI=2.*PI
  U=4.*PI*1.e-7 ! magn perm of free space
  rnu=0. !wavenumber=0,plane wave model

  do I=1, npoint
    F(i)=FLOAT(i)*60.
    OMG=TOPI*F(i)
  call layered_Earth(sigmaON2,dON2,OMG,rnu,NON,cON2)
C   Z(I)=(0.,1.)*TOPI*F(i)*cQ(1) ! in m/s
   Z1(i)=(0.,1.)*TOPI*F(i)*cON2(1)*U !in Ohm
   ZR(i)=REAL(Z1(i))
   ZI(i)=AIMAG(Z1(i)) ! not inverted
   ZZ(i)=CMPLX(ZR(i),ZI(I))
   ZA(i)=SQRT(ZR(i)*ZR(i)+ZI(i)*ZI(i))
   Zph(i)=ATAN(ZI(i)/ZR(i))*180./PI
   P(i)=ZZ(i)/(0.,1.)/OMG/U

```

C Calculations of the E of each phase normalised by currents/change to function

```

  EaL(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HaL+2*P(i))**2+XaL**2)/
  *SQRT(HaL**2+XaL**2))
  EaR(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HaR+2*P(i))**2+XaR**2)/
  *SQRT(HaR**2+XaR**2))
  EbL(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HbL+2*P(i))**2+XbL**2)/
  *SQRT(HbL**2+XbL**2))
  EbR(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HbR+2*P(i))**2+XbR**2)/
  *SQRT(HbR**2+XbR**2))
  EcL(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HcL+2*P(i))**2+XcL**2)/
  *SQRT(HcL**2+XcL**2))
  EcR(i)=-((0.,1.)*OMG*U/topi)*clog(CSQRT((HcR+2*P(i))**2+XcR**2)/
  *SQRT(HcR**2+XcR**2))

```

end do !end of I

C calculations of total electric field for each harmonic

```

  E(1)=Cur(1)*((EaL(1)+EaR(1))+(EbL(1)+EbR(1))*cos(240.*PI/180.)+
  * (EcL(1)+EcR(1))*cos(120.*PI/180.))
  E(2)=Cur(2)*((EaL(2)+EaR(2))+(EbL(2)+EbR(2))*cos(120.*PI/180.)+

```

```

* (EcL(2)+EcR(2))*cos(240.*PI/180.)
  E(3)=Cur(3)*((EaL(3)+EaR(3))+(EbL(3)+EbR(3))+
* (EcL(3)+EcR(3)))

  open(UNIT=20,file='Efield_Harmonics')
  write (20,*) 'Harmonics_E_Field'
  write(20,*) 'Hz skin-depth Za Zph Eamp Ephase'

  do i=1,npoint
    K=i*60
    AMPLE(i)=SQRT(REAL(E(i))**2+AIMAG(E(i))**2)
    PhasE(i)=ATAN(AIMAG(E(i))/REAL(E(i)))*180./PI

  write (20,20) k,P(i),Za(i),Zph(i),REAL(E(i)),AIMAG(E(i))
20  Format (I3,6(2X,F17.10))
  end do
  close(UNIT=20)

C222 continue
  stop
  end

  subroutine layered_Earth(sigma,d,omega,rnu,N,c)
! input parameters
!   sigma(N) - array of layer's conductivity
!   d(N) - layers depth
!   omega - external frequency
!   rnu - wavenumber
!   N - number of layers ( dimension of arrays)
! output parameters
!   c1 - skin-depth of the 1-st layer
  real d(N),sigma(N)
  complex c(n),alpha2,gamma2,R,gamma
  real pi/3.1415926/
  Rmu0=4.0*pi*1E-7
  gamma2=Rnu*Rnu+(0.,1.)*omega*Rmu0*Sigma(n)
  C(N)=1./SQRT(gamma2)

  do i=N-1,1,-1
    Alpha2=omega*Rmu0*Sigma(i)
    gamma2=Rnu*Rnu+(0.,1.)*Alpha2
    gamma=SQRT(gamma2)
    R=(1.-gamma*C(i+1))/(1.+gamma*C(i+1))
    C(i)=(1.-R*exp(-2.*gamma*d(i)))/gamma/(1.+R*exp(-2.*gamma*d(i)))
  end do

  return
  end

```

## 7.2 FORTRAN Code 2

PROGRAM TO CALCULATE PSP induced by harmonics  
with frequency-dependent pipeline impedances  
electric field in the ground induced by harmonics used as the input

```
parameter is number of harmonics
parameter (nharm=3)
parameter (Nmax=3000) !1.E+4
real F(nharm),DatEx(nmax,nharm),
& DataEy(nmax,nharm),PSP(nmax,nharm), Cur(nmax,nharm)
COMPLEX Ex(nharm),Ey(nharm),Z,Y, DatZ(nharm),DatY(nharm)
character*80 fileIN1
complex ZI(2),V,CI,EAX,EAY, DatP(nharm), DatCurr(nharm)

real sect(2), fi(2), xx(2)
common /sections/ nsection, sect, fi,EAX,EAY

PI=acos(-1.0)
topi=2.*PI
C dt=1.E-4 !sampling rate in sec.
Tmin=1./(60.*nharm)
dt=Tmin/1000. ! at least 1000 points in the smallest period

C Nmax=180./(nharm*F*dt) ! maximum size of array for waveform calculations

C electric field for Ontario case
C EX(1)=1000.*(-3.6E-7,4.9E-3) ! V/km
C EX(2)=1000.*(-1.25E-8,9.78E-5) !V/km
C EX(3)=1000.*(-1.08E-2,-4.87E-2) !V/km

C Electric field for Ontario with top water
EX(1)=1000.*(-0.0000102208,0.0049109999) ! V/km
EX(2)=1000.*(0.0000000058,0.0000985511) !V/km
EX(3)=1000.*(-0.0083882259,-0.0189491455) !V/km

C CALCULATIONS OF PSP

C pipeline geometry

nsection=2 !5

sect(1)= 250. !115.
sect(2)=200. !55.
C sect(3)=30.
C sect(4)=90.
```

```

C      sect(5)=161.

      fi(1)=1. !3.
      fi(2)=3. ! 18.
C      fi(3)=55.
C      fi(4)=3.
C      fi(5)=25.

      xx(1)=0.0
      xx(2)=xx(1)
      do i=1,nsection
      xx(2)=xx(2)+sect(i)
      end do

```

```

C pipeline terminating impedances
      ZI(1)= (0.1,0.0) !(100.0,0.)
      ZI(2)= (0.1,0.0) !(100.0,0.)
      Np=901
      x=185. !coordinate for PSP calculation

```

C PSP calculations

```

      Do i=1,nharm
      F(i)=60.*i
      EAX=EX(i)
      EAY=(0.,0.) !EY(i) line parallel to pipeline
      OMG=topi*F(i)
      call TLINE(OMG,x,XX,EAX,EAY,ZI,Np,V,CI,Z,Y)
      DatP(i)=V
      DatCurr(i)=CI
      DatZ(i)=Z
      DatY(i)=Y
      end do

```

C Check the reading

```

C      open(UNIT=35,file='StartE.dat')
C      do i=1,nharm
C      write(35,35) f(i), EX(i),DatP(i),DatCurr(i)
C 35  Format(7(2X,E15.7))
C      end do
C      goto 110

```

C conversion of E to time domain (for three fruequency only)

```

      do i=1,nharm
      F(i)=60.*i

```

```

DO k=1,Nmax !number of periods for waveform calculations
!dt=1.E-4 is sampling rate for waveform calculations

```

```

DATEX(k,i)=SQRT((REAL(EX(i))**2+(AIMAG(EX(i))**2)*
&sin(topi*F(i)*dt*k+atan(AIMAG(Ex(i))/REAL(Ex(i))))

```

```

C DATAEY(k,i)=SQRT(REAL(EY(i)**2)+AIMAG(EY(i)**2))*
C &sin(topi*F(i)*dt*k+atan(AIMAG(EY(i))/REAL(EY(i))))

```

```

PSP(k,i)=SQRT((REAL(DatP(i))**2+(AIMAG(DatP(i))**2)*
&sin(topi*F(i)*dt*k+atan(AIMAG(DatP(i))/REAL(DatP(i))))

```

```

CUR(k,i)=SQRT((REAL(DatCurr(i))**2+(AIMAG(DatCurr(i))**2)*
&sin(topi*F(i)*dt*k+atan(AIMAG(DatCurr(i))/REAL(DatCurr(i))))

```

```

end do

```

```

end do

```

```

open(UNIT=36,file='Pipeline_Impedances.dat')
do i=1,nharm
F(i)=60.*i

```

```

write(36,36) F(i), DATZ(i),DATY(i)
36 Format(5(2X,E15.7))
end do

```

```

C put results in file
open(UNIT=41,file='Harmonics_psp.dat')
write (41,*) 'Harmonics_PSP_Current'

```

```

C output
write(41,*) ' sec, dt, F(Hz), Ex (mV/km), PSP(mV), Current(A)'
do i=1,nharm
DO k=1,nmax
write (41,41) k,F(i),DATEX(k,i),PSP(k,i)
&,Cur(k,i)
41 Format (I5,4(2X,E15.7))

```

```

end DO

```

```

end do

```

```

110 continue

```

```

stop
end

```

subroutine TLINE(OMG,x,XX,EXA,EYA,ZI,Np,V,CI,Z,Y)

c INPUT PARAMETERS:

c x is coordinate along pipeline where V and CI have to be computed

c XX is a real array of 2 elements: XX(1) beginnig of pipe, XX(2) end of  
C pipe

c  $Z=R+jwL$  - complex impedance per unit length;

c  $Y=G+jwC$  - complex admittance per unit length;

c ZI - complex array of 2 elements: Z(1) terminating impedance at XX(1)

c                                   Z(2) terminating impedance at XX(2)

c Np - number of points for the integration

c OUTPUT PARAMETERS:

c V - voltage at x

c CI - current at x

complex Z,Y,ZI(2),V,CI

complex WNs,ZZ,Zint,Zext,eax,eay

real x, XX(2),sect(2),fi(2),XXX(2)

complex Z0,GAMMA,R1,R2,EXA,EYA

complex K1,K2

complex P, Q ! complex integral-function

complex CNUM1,CNUM2,CDEN

COMMON /PARAM/ Z0,XXX,GAMMA,N

common /sections/ nsection, sect, fi,eax,EAY

C Formulas for  $Z=int+ext$  and  $Y=G+i*OMG*C$

C PIPELINE INPUT VALUES:

C Rext-external radius of steel,

C Rint-internal

C D-coating thickness

C Cp-conductivity of steelpipe

C Cs-conductivity of soil

C Cc-conductivity of coating

C EPSc-permittvity of the coating

C Up-magnetic permeability of the steel

C Us-magnetic permeability of the soil

C OMG-frequency

PI=acos(-1.0)

topi=2.\*PI

XXX(1)=XX(1)

XXX(2)=XX(2)

C goto 1



```

Rext=0.457
Rint=0.448
D=0.00025 !
Cp=5.E+6 !
Cs=1.E-3
Cc=20.E-6 !coat. conductance per meter^2 change to 20.
EPSc=3./(18.E+9*topi) ! change to 3.
Up=150.*(2.*TOPI)/1.E+7 !change to 150
Us=(2.*TOPI)/1.E+7
Thk=Rext-Rint
RR=1./(pi*Cp*(Rext**2-Rint**2))
Rav=(Rext+Rint)/2.
WNs=CSQRT((0.0,1.0)*OMG*Us*Cs) !complex
ZZ=RR*CSQRT((0.0,1.0)*Up*OMG/topi/thk/RR/Rav) !complex
Zint=ZZ*thk*(CEXP(ZZ/RR*thk)+CEXP(-ZZ/RR*thk))/(CEXP(ZZ/RR*thk)
& -CEXP(-ZZ/RR*thk)) !complex
Zext=(0.0,1.0)*OMG*Us/topi*CLOG(2./WNs*Rext) !complex
Z=1000.*(Zint+Zext) !Ohm/km, complex
G=topi*Cc*(Rext+d)*1.E+3 !S/km
C=topi*EPSc/(LOG(1.+d/Rext))*1.E+3 !F/km
Y=(G+(0.0,1.0)*OMG*C) !S/km,complex

```

1 continue

C Z=(0.00782,0.00782)

C Y=(0.00576,0.0)

N=Np

c define some constants

```

Z0=CSQRT(Z/Y) ! units [ohms]
GAMMA=CSQRT(Z*Y) ! units [length^-1]
R1=(ZI(1)-Z0)/(ZI(1)+Z0) ! reflection coef at XX(1)
R2=(ZI(2)-Z0)/(ZI(2)+Z0) ! reflection coef at XX(2)
DX=XX(2)-XX(1)

```

```

CNUM1=R2*P(XX(2))*CEXP(-GAMMA*XX(2))
&-Q(XX(1))*CEXP(GAMMA*XX(2))

```

```

CDEN=CEXP(GAMMA*DX)-R1*R2*CEXP(-GAMMA*DX)

```

```

K1=R1*CEXP(GAMMA*XX(1))*CNUM1/CDEN

```

```

CNUM2=R1*Q(XX(1))*CEXP(GAMMA*XX(1))
&-P(XX(2))*CEXP(-GAMMA*XX(1))

```

c denominator is the same CDEN

```

K2=R2*CEXP(-GAMMA*XX(2))*CNUM2/CDEN

```

```

CI=(K1+P(x))*CEXP(-GAMMA*x)
&+(K2+Q(x))*CEXP(GAMMA*x)
V=Z0*((K1+P(x))*CEXP(-GAMMA*x)
&-(K2+Q(x))*CEXP(GAMMA*x))

```

```

return
end

```

```

FUNCTION P(x)
COMMON /PARAM/ Z0,XX,GAMMA,N
    common /sections/ nsection, sect, fi,EAX,EAY
COMPLEX P,PP,Z0,GAMMA,Epipе,EAX,EAY
REAL x,XX(2),sect(2),fi(2)

```

```

if(mod(N,2).eq.0) then
    write(*,*) 'mod(N,2)=0 in P'
    stop ''
end if

```

```

if(n.le.2) stop ' P: N<=2'

```

```

PP=CEXP(GAMMA*XX(1))*Epipе(XX(1))
PP=PP+ CEXP(GAMMA*x)*Epipе(x)

```

```

h=(x-XX(1))/(N-1.)

```

```

do i=2,N,2
    xt=XX(1)+h*(i-1)
    PP=PP+ (4.0,0.)*CEXP(GAMMA*xt)*Epipе(xt)
end do

```

```

do i=3,N-1,2
    xt=XX(1)+h*(i-1)
    PP=PP+ (2.0,0.)*CEXP(GAMMA*xt)*Epipе(xt)
end do

```

```

P=1./(2.,0.)/Z0*PP*h/3.

```

```

RETURN
END

```

```

FUNCTION Q(x)
COMMON /PARAM/ Z0,XX,GAMMA,N
    common /sections/ nsection, sect, fi,EAX,EAY
COMPLEX Q,QQ,Z0, GAMMA, Epipе,EAX,EAY
REAL x,XX(2),sect(2),fi(2)

```

```

if(mod(N,2).eq.0) then
  write(*,*) 'mod(N,2)=0 in Q'
  stop ''
end if

if(n.le.2) stop ' Q: N<=2'

QQ=CEXP(-GAMMA*XX(2))*Epipe(XX(2))
QQ=QQ+ CEXP(-GAMMA*x)*Epipe(x)

h=(XX(2)-x)/(N-1.)

do i=2,N,2
  xt=x+h*(i-1)
  QQ=QQ+(4.0,0.)*CEXP(-GAMMA*xt)*Epipe(xt)
end do

do i=3,N-1,2
  xt=x+h*(i-1)
  QQ=QQ+ (2.0,0.)*CEXP(-GAMMA*xt)*Epipe(xt)
end do

Q=1./(2.0,0.)/Z0*QQ*h/3.

RETURN
END

function Epipe(x)
  common /sections/ nsection, sect, fi,EAX,EAY
  complex Epipe,EAX,EAY
  real x,sect(2), fi(2)

  RG=3.1415926/180.
  SECTA=SECT(1)
  Epipe=EAY*sin(fi(1)*RG)+EAX*cos(fi(1)*RG)
  Do i=2,nsection
  IF((X-secta).GT.0.0) Epipe=EAY*sin(fi(i)*RG)+EAX*cos(fi(i)*RG)

  SECTA=SECTA+sect(i)
  end do
return
end

```

## 8. References

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